

From 2D to 1D in β -naphthyne: A porous carbon allotrope merging graphyne and naphthylene

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ABSTRACT

Two-dimensional (2D) carbon-based materials have attracted considerable interest due to their diverse structural and electronic properties, making them ideal for next-generation flat electronics. Among these materials, metallic-like porous structures offer advantages such as tunable charge transport and high surface area, which are essential for energy storage applications. In this study, we introduce β -naphthyne, a novel 2D carbon allotrope composed of naphthyl units interconnected by octagonal rings. First-principles calculations confirm its dynamic and thermal stability, demonstrating its theoretical feasibility. Furthermore, we demonstrate that Young's modulus ranges from 43.71 N/m to 165.88 N/m, indicating an anisotropic mechanical response. Optical analysis reveals absorption activity in the infrared (IR) and ultraviolet (UV) regions. The derived 1D structures were also analyzed, revealing a Dirac cone and a transition from metallic to semiconducting behavior. These findings establish β -naphthyne as a promising material for energy storage and optoelectronic technologies.

1. Introduction

Two-dimensional (2D) carbon-based materials have garnered significant attention due to their mechanical, electronic, and optical properties, making them promising candidates for energy storage and nanoelectronics applications [1–4]. Among them, porous 2D carbon allotropes offer additional advantages while maintaining a high surface-to-volume ratio, exhibiting tunable electronic structures and enhanced ion diffusion, which are crucial for hydrogen storage, supercapacitors, and metal-ion batteries [5–9]. Incorporation of ordered nanopores into carbon frameworks can introduce different electronic properties, such as Dirac cone semimetallicity, which pave the way for the development of high-performance energy materials [10–12].

Several porous carbon allotropes have been synthesized or theoretically predicted in recent years, each exhibiting distinct structural and electronic characteristics [13–17]. For instance, graphdiyne, a synthesized material featuring acetylenic linkages, exhibits enhanced electrochemical performance due to its high charge carrier mobility and tunable band gap [18]. Other experimentally synthesized materials include the biphenylene network, which was recently fabricated via on-surface synthesis and displays metallic behavior, holding potential

for nanoelectronic applications [19]. In contrast, several computationally predicted structures [20–28] have been characterized, revealing distinct structural, electronic, optical, and mechanical properties. However, the experimental realization of some of these structures remains a challenge [29]. However, theoretical studies are essential for guiding experimental synthesis by predicting stability, mechanical response, and electronic characteristics, enabling researchers to design novel materials with desirable functionalities before fabrication [30,31].

Building upon recent advances in 2D carbon-based materials for surface electronics, we introduce the graphyne-like β -naphthylene (β -naphthyne) as a novel porous 2D carbon allotrope characterized by naphthalene molecules covalently linked through acetylenic groups at all their dehydrogenated terminations, forming a rectangular symmetry. β -naphthyne is a graphyne-like version of β -naphthylene, previously proposed by Paz et al. [32] in which the allotrope is predicted from the naphthalene precursor in its so-called β phase, configuring a 2D carbon allotrope with 4-, 6-, and 10-membered rings. In the structure proposed here, the linkage between naphthalene molecules is replaced by acetylenic groups, transforming the originally sp^2 -bonded system into a hybrid sp - sp^2 configuration, now exhibiting 6- and

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8-membered rings, along with nanopores consisting of 18 carbon-membered units.

Regarding the challenges of synthesizing such materials, naphthalene units no longer serve as a precursor for a potential synthesis of β -naphthylene. Instead, octamethylnaphthalene ($C_{18}H_{24}$), synthesized in 1972 by Hart and Oku [33], emerges as a viable structural building block for bottom-up synthesis strategies, such as surface-assisted polymerization or templated chemical vapor deposition (CVD), providing a feasible route for future syntheses of β -naphthylene [34–36].

This study presents the novel β -naphthylene structure, investigated through first-principles calculations. Using density functional theory (DFT), we explored its electronic properties, including band structures, electron localization function (ELF), and Bader charge distribution. The thermal and dynamic stabilities were verified using *ab initio* molecular dynamics (AIMD) and phonon dispersion, respectively. A comprehensive mechanical evaluation was conducted by analyzing elastic constants, with stability confirmed via the Born-Huang criteria. In addition, the effect of 1D confinement was evaluated by analysis of β -naphthylene nanoribbons. The findings reported here not only highlight the β -naphthylene as a 2D reliable carbon allotrope but also suggest its potential as a platform for diverse applications, such as energy storage, gas detection, and electronics.

2. Methodology

To explore the structural, electronic, mechanical, and optical properties of β -naphthylene, we conducted first-principles simulations based on DFT [37]. The computational framework was implemented using the Vienna Ab Initio Simulation Package (VASP) [38]. For electron exchange and correlation effects, we employed the generalized gradient approximation (GGA) following the Perdew–Burke–Ernzerhof (PBE) functional [39]. The interaction between valence and core electrons was treated using the projector-augmented wave (PAW) method [40].

The kinetic energy cutoff for the plane-wave basis set was set to 520 eV, ensuring an accurate representation of electronic states. The structural optimizations were performed using the conjugate gradient method until the force on each atom was lower than 0.01 eV/Å, and the total energy change between successive steps was below 10^{-5} eV. A vacuum spacing of 15 Å along the out-of-plane direction was introduced to prevent interactions between periodic images. The Brillouin zone was sampled using a Γ -centered Monkhorst–Pack k -point mesh of $3 \times 3 \times 1$ ($18 \times 18 \times 1$) for monolayers and $3 \times 1 \times 1$ ($18 \times 1 \times 1$) for nanoribbons for optimization (density of states) calculations. Dispersion interactions were included using Grimme's DFT-D2 method [41], accounting for van der Waals forces to accurately describe interatomic interactions in porous carbon frameworks.

To evaluate the dynamical stability of β -naphthylene, phonon dispersion calculations were performed using density functional perturbation theory (DFPT) as implemented in the Phonopy package [42]. Additionally, *ab initio* molecular dynamics (AIMD) simulations were conducted at 300 K for 5 ps in the NVT ensemble, employing the Nosé–Hoover thermostat [43,44] to assess its thermal robustness.

The optical properties were investigated through calculations of the frequency-dependent dielectric function, from which absorption coefficients and reflectance spectra were derived [45], providing insight into the potential optoelectronic applications of β -naphthylene.

The energetic stability of β -naphthylene was analyzed by its cohesive energy (E_{coh}):

$$E_{\text{coh}} = \frac{E_{\beta\text{-naphthylene}} - \sum_i n_i E_i}{\sum_i n_i}, \quad (1)$$

where $E_{\beta\text{-naphthylene}}$ is the total energy of the structure, E_i represents the energy of isolated C atoms, and n_i is the atom count of each element. The same approach was used to obtain the E_{coh} for other 2D structures discussed in this work.

To analyze the vibrational properties of β -naphthylene, the coupled perturbed HF/Kohn–Sham algorithm [46] implemented in the computational package CRYSTAL17 [47] was used. The calculations in CRYSTAL17 were performed using triple-zeta valence with polarization (TZVP) basis set [48] together with the B3LYP functional. The structure was optimized by checking the root mean square (RMS) and the absolute value of the largest component of both the gradients and the estimated displacements. The convergence criteria used in the optimization for RMS and the largest component for gradient (0.00030 a.u. and 0.00045 a.u.) and displacement (0.00120 a.u. and 0.00180 a.u.), respectively. The reciprocal space was sampled using Pack–Monkhorst and Gilat nets with sublattice and a shrinking factor of 12.

3. Results

3.1. Structure and stability

β -naphthylene is characterized by a rectangular conventional unit cell that belongs to the Cmm (no. 65) space group, with lattice parameters $a = 10.74$ Å and $b = 11.72$ Å (see Fig. 1a). The primitive cell comprises lattice parameters $a = b = 7.82$ Å, $\alpha = \beta = 90^\circ$, and $\gamma = 93.17^\circ$. The conventional unit cell consists of five crystallographically distinct carbon atoms located in C1 (0.764, 0.378, 0.500), C2 (0.359, 0.181, 0.500), C3 (0.566, 0.282, 0.500), C4 (0.367, 0.392, 0.500), and C5 (0.932, 0.000, 0.500). This structure forms a 2D carbon framework comprising 6–8 carbon-membered rings, where octagonal rings interconnect naphthyl fragments. For comparison purposes, Table 1 details the lattice parameters (a , b , $\alpha = \beta$, and γ), space group (SPG), band gap energy (E_{gap}), and cohesive energy (E_{coh}) for several carbon monolayers. The comparison shows that the cohesive energy (E_{coh}) of β -naphthylene (−7.25 eV/atom) is competitive with other 2D carbon allotropes, being close to that of T-graphene (−7.45 eV/atom) and graphenylene (−7.33 eV/atom), but lower than that of pure graphene (−8.02 eV/atom). The E_{coh} reflects the energy required to disassemble a solid into isolated atoms, serving as a key indicator of the structural stability of the material. Therefore, while β -naphthylene shows slightly weaker cohesion than graphene, its value remains sufficiently high to imply a stable 2D structure.

To evaluate the dynamical stability of β -naphthylene, the phonon dispersion was computed along the high-symmetry paths of the Brillouin zone, as depicted in Fig. 1b. The dispersion confirms the absence of imaginary frequencies, indicating that the material is dynamically stable. At the Γ -point, three acoustic modes with a characteristic linear behavior. Several phonon branches remain relatively flat in the 18–25 THz range, suggesting low group velocity and localized vibrational modes. In contrast, the bands display quadratic dispersion in the 25–35 THz range, indicating high phononic mobility. Finally, vibrational modes with flat characteristics above 60 THz correspond to the sp-hybridized bonds in β -naphthylene, a common feature in graphyne-like systems [55].

One critical aspect in assessing the viability of novel 2D materials is their thermal stability under ambient conditions. In this way, AIMD simulations were performed at 300 K for a total duration time of 4 ps. Fig. 2a depicts the evolution of the total energy fluctuations during the simulation, along with the final atomic configuration (Fig. 2b).

The total energy stabilizes after the first picosecond of the simulation. After this phase, energy fluctuations remain minimal, consistently within a range of less than 1 eV, indicating structural robustness. The absence of significant deviations suggests that no phase transitions or reconstruction events occur during the simulation, further supporting the stability of the material at room temperature. The final atomic configuration, presented in Fig. 2b, further corroborates this observation. The top and side views reveal that the structural integrity of β -naphthylene is preserved, with only minor out-of-plane distortions emerging.

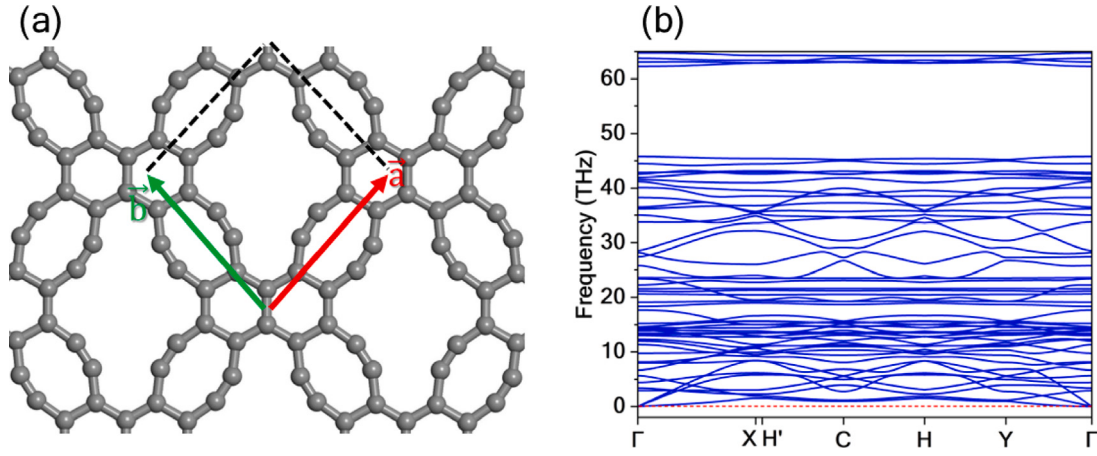


Fig. 1. (a) Top view of the optimized β -naphthylene unit cell (primitive lattice vectors $a = b = 7.82 \text{ \AA}$); (b) phonon-dispersion along the high-symmetry directions of the Brillouin zone.

Table 1

Lattice parameters (a and b in $\text{\AA}$, $\alpha = \beta$, and γ in degrees), space group (SPG), band gap energy (E_{gap}) in eV, and cohesive energy (E_{coh}) in eV/atom obtained for β -naphthylene and other relevant 2D carbon allotropes at DFT/GGA-PBE level. The values reported in this table were calculated in this study.

	a	b	$\alpha = \beta$	γ	SPG	E_{gap}	E_{coh}
β -naphthylene (this work)	10.74	11.43	90	90	Cmm (65)	Metallic	-7.25
T-graphene [49]	3.45	3.45	90	90	$P4/mmm$ (123)	Metallic	-7.45
Twin-graphene [50]	6.14	6.14	90	120	$P6/mmm$ (191)	0.73	-7.08
PHE-graphene [51]	5.73	5.73	90	120	$P6m2$ (187)	Metallic	-7.56
Penta-graphene [21]	3.64	3.64	90	90	$P4_2/m$ (113)	2.19	-7.13
Graphyne [52]	9.46	9.46	90	120	$P6/mmm$ (191)	Metallic	-7.20
Graphenyldiene [28]	9.32	9.32	90	120	$P6/mmm$ (191)	0.78	-6.92
Graphenylene [27,53]	6.77	6.77	90	120	$P6/mmm$ (191)	0.04	-7.33
Biphenylene [19]	3.80	4.50	90	90	$Pmmm$ (47)	Metallic	-7.47
Graphene [54]	2.47	2.47	90	120	$P6/mmm$ (191)	Metallic	-8.02

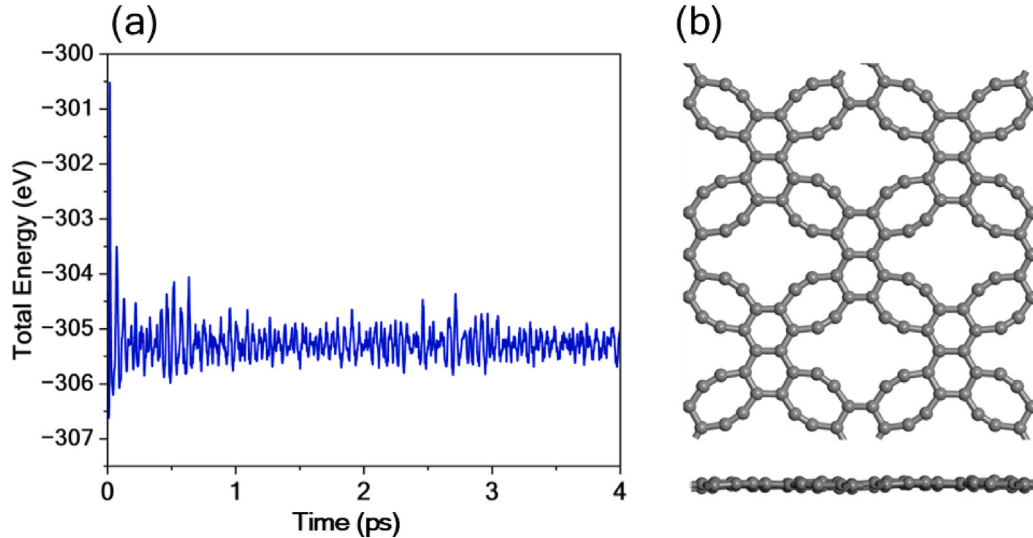


Fig. 2. (a) Evolution of the total energy at 300 K over a 4 ps AIMD run; (b) final atomic configuration (top and side views) after the simulation.

3.2. Electronic analysis

We performed a detailed analysis of the band structure to investigate the electronic properties of β -naphthylene. The results, depicted in Fig. 3, confirm that β -naphthylene exhibits metallic behavior with a band partially occupied crossing the Fermi level (E_F) closer to the Γ point. The states that compose this band can act as acceptor states, receiving the electronic density, which can be useful for an effective interaction with metal adatoms. In the valence band (VB), one can observe that

the bands exhibit high dispersion, with several crossings, indicating a strong coupling between the electronic states. On the other hand, when the conduction band (CB) is analyzed, the lowest energy band touches E_F two times at the points H' and H . This band presents a high dispersion, around 2 eV.

The projected density of states (PDOS) analysis provides a more detailed understanding of the electronic structure, shown in Fig. 4, which decomposes the total density of states into atomic orbital contributions. The VB and CB predominantly comprise p orbitals, with the p_y states

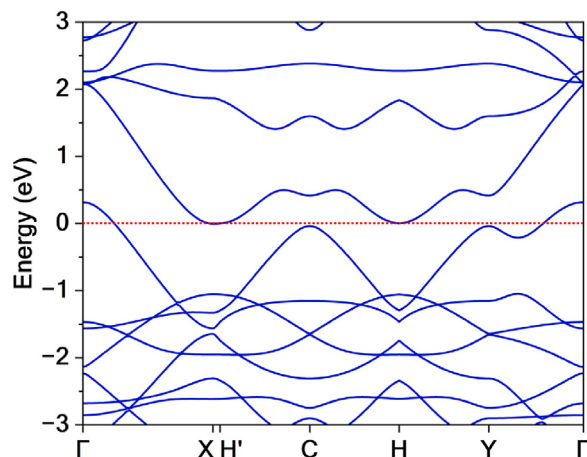


Fig. 3. Electronic band structure of β -naphthylene along the principal Brillouin-zone paths; the red dashed line marks the Fermi level (E_F) (set to 0 eV).

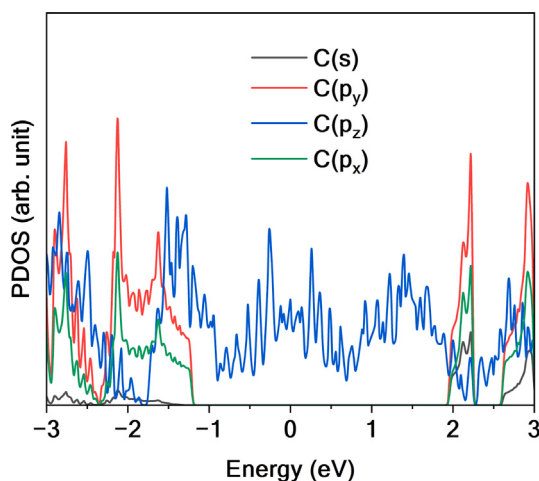


Fig. 4. Orbital-resolved projected density of states (PDOS) for β -naphthylene.

contributing most significantly, closely followed by the p_x states. The similar energy distribution profiles of the p_y and p_x orbitals suggest a strong orbital coupling in these energy regions. However, near the E_F , it becomes evident that only the p_z orbitals contribute, indicating a pronounced π -character in this energy range.

To further investigate the electronic properties of β -naphthylene, we analyzed the spatial distribution of its frontier crystalline orbitals, specifically the highest occupied crystalline orbital (HOCO) and the lowest unoccupied crystalline orbital (LUCO), as shown in Fig. 5. Fig. 5a illustrates HOCO, which is predominantly composed of π -type orbitals located on the naphthyl units. The electron density is highly delocalized across the hexagonal rings, indicating a strong conjugation network. The side view further confirms the planarity of the electronic distribution, reinforcing the 2D nature of the conduction pathways.

Fig. 5b shows the LUCO, which also exhibits a π -character, but with electron density more prominently localized around the octagonal rings. This spatial separation between the HOCO and LUCO suggests a moderate charge transfer character, which could influence optical excitations and recombination rates. The LUCO distribution indicates that electron transport might be primarily facilitated through the octagonal framework, complementing the hole conduction pathways of the HOCO.

The distinct spatial location of HOCO and LUCO further supports the metallic nature of β -naphthylene, as both orbitals maintain a significant

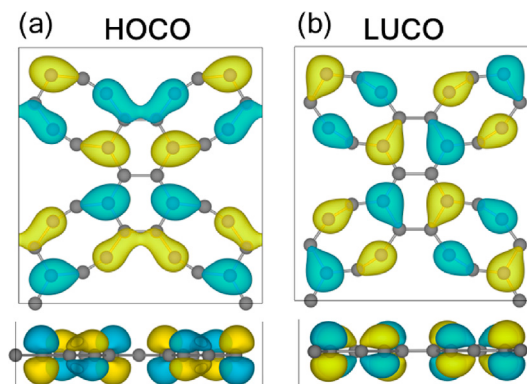


Fig. 5. Real-space distribution of (a) the highest occupied crystalline orbital (HOCO) and (b) the lowest unoccupied crystalline orbital (LUCO) of β -naphthylene.

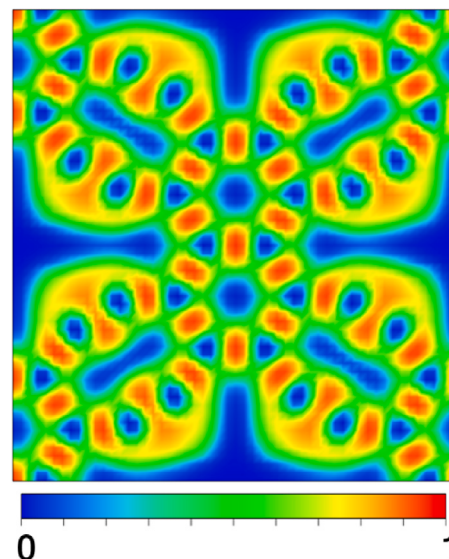


Fig. 6. Electron localization function (ELF) of β -naphthylene.

density near the E_F . The overlap of these orbitals could also contribute to strong optical absorption, suggesting potential applications in photodetectors and optoelectronic devices, as will be discussed later.

To gain a deeper understanding of the charge distribution and bonding nature of β -naphthylene, we analyzed its electron localization function (ELF), as shown in Fig. 6. The ELF provides valuable information on the degree of electron localization, where a value of 1 indicates strongly localized electrons (e.g., covalent bonds or lone pairs), a value of 0.5 corresponds to a homogeneous electron gas-like distribution, and a value of 0 represents regions with minimal or no electron density.

In β -naphthylene, the naphthyl fragments exhibit high ELF values, as expected due to the sp^2 hybridization. In contrast, the octagonal rings exhibit lower ELF values, indicating greater electron delocalization than the naphthyl units. This behavior can be attributed to sp -hybridized bonds, which result in a more diffuse electron density along the in-plane bonds.

The ELF map further supports the conclusion that β -naphthylene exhibits a highly delocalized electronic structure, with distinct regions contributing to different transport properties. The bond interactions within the naphthyl and octagonal rings, i.e., the combination of sp and sp^2 hybridizations, create a unique electronic environment, making this material a promising candidate for nanoelectronic and optoelectronic applications, where controlled charge delocalization and transport are essential.

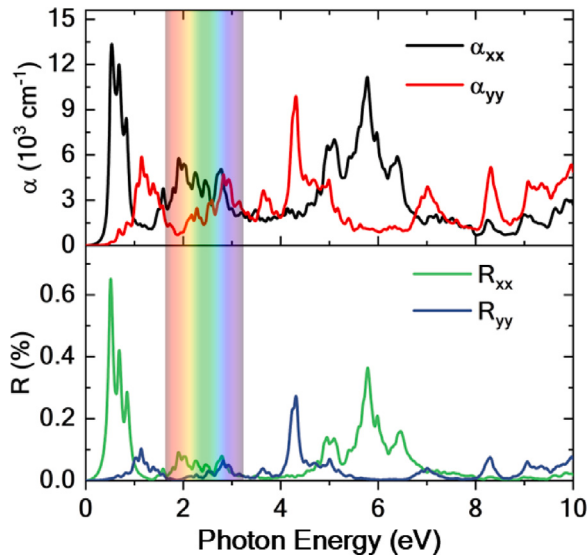


Fig. 7. Optical response of β -naphthylene: (top) absorption coefficient (α) for light polarized along x (black) and y (red); (bottom) reflectance (R) for the same polarization directions, x (green) and y (blue).

3.3. Optical properties

To address the optical properties of β -naphthylene, we calculated its absorption coefficient (α) and reflectance (R), as illustrated in Fig. 7. The visible light range is highlighted in the figure to facilitate the identification of its transparency and absorption characteristics.

β -Naphthylene exhibits strong absorption in the IR region, with the absorption coefficient reaching approximately $14 \times 10^3 \text{ cm}^{-1}$ for the x component. This result suggests that the material could be well-suited for IR photodetectors and thermal shielding applications. Additionally, the y component shows noticeable absorption in the near-IR range, further reinforcing its potential for IR-active optoelectronic devices.

The absorption onset for the x -component is observed near 0.4 eV, which corresponds to a wavelength of approximately $3.1 \mu\text{m}$. This places it at the boundary between the short-wave infrared (SWIR, $1.4\text{--}3 \mu\text{m}$) and the mid-wave infrared (MWIR, $3\text{--}5 \mu\text{m}$) regions. These spectral ranges are relevant for various commercial IR devices such as gas sensors [56] and environmental monitoring [57] systems, indicating that β -naphthylene could be a promising candidate for MWIR photodetection technologies.

The material displays relatively low absorption in the visible spectrum, with the maximum absorption coefficient reaching $6 \times 10^3 \text{ cm}^{-1}$ between the red and orange spectral regions for the x component. This trend suggests that β -naphthylene is mainly transparent to visible light, making it a promising candidate for transparent conducting films or low-energy-loss optical coatings. The reflectance in this range remains minimal. β -naphthylene exhibits strong absorption peaks in the UV region, with significant transitions occurring around 4.5 eV (y component) and 6 eV (x component).

3.4. Vibrational properties

Vibrational spectroscopy provides essential insights into the bonding nature and structural integrity of materials. To assess the phonon behavior of β -naphthylene, we computed both the Raman and IR spectra, as shown in Fig. 8. These spectra allow us to distinguish between Raman-active and IR-active vibrational modes, shedding light on the symmetry and electronic environment of atomic vibrations in the material.

The vibrational modes of β -naphthylene can be represented by $\Gamma_{\text{vib}} = 5B_{1u} + 9B_{2u} + 9B_{3u} + 4B_{2g} + 4A_u + 5B_{3g} + 9B_{1g} + 9A_g$. Between these

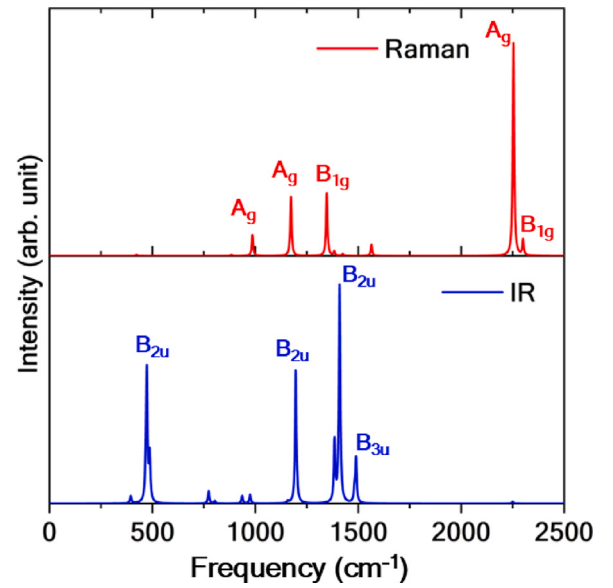


Fig. 8. Calculated vibrational spectra of β -naphthylene: (top) Raman spectrum and (bottom) infrared spectrum.

modes, three are acoustic ($B_{1u} + B_{2u} + B_{3u}$), four are silent ($4A_u$), 23 are IR-active ($5B_{1u} + 9B_{2u} + 9B_{3u}$), and 27 are Raman-active ($4B_{2g} + 5B_{3g} + 9B_{1g} + 9A_g$).

Analyzing the Raman spectrum (Fig. 8, top panel) reveals several distinct peaks, with the most intense modes A_g and B_{1g} occurring at low and intermediate frequencies. Four prominent bands are visible at 985 cm^{-1} (A_g), 1172 cm^{-1} (A_g), 1383 cm^{-1} (B_{1g}), and 2253 cm^{-1} (A_g). Strong A_g symmetric modes suggest collective in-plane stretching and breathing vibrations of the carbon framework.

However, by verifying the IR spectrum (Fig. 8, bottom panel), several intense peaks associated with asymmetric stretching and bending modes are shown. Four intense bands are noticed at 472 cm^{-1} (B_{2u}), 1196 cm^{-1} (B_{2u}), 1409 cm^{-1} (B_{2u}), and 1489 cm^{-1} (B_{3u}) are observed. The B_{2u} and B_{3u} indicate dipole-active vibrations arising from charge asymmetry within the structure. The strongest IR peak, located around 1500 cm^{-1} , is related to C–C stretching motions similar to those observed in other 2D carbon-based materials [58,59]. The presence of multiple low-frequency peaks suggests the existence of out-of-plane bending and deformation modes.

3.5. Mechanical properties

The mechanical properties of β -naphthylene were assessed by calculating its elastic constants, which are crucial to determining the material's structural stability and mechanical response. The obtained values for the independent elastic constants are $C_{11} = 169.10 \text{ N/m}$, $C_{22} = 173.64 \text{ N/m}$, $C_{12} = 42.68 \text{ N/m}$, $C_{16} = -22.32 \text{ N/m}$, $C_{26} = -17.49 \text{ N/m}$, and $C_{66} = 20.98 \text{ N/m}$. According to the Born-Huang stability criteria [60], the mechanical stability of 2D oblique systems is ensured if $C_{11} > 0$, $C_{11}C_{22} > C_{12}^2$, and $\det(C_{ij}) > 0$. Substituting the calculated values, these conditions are satisfied. Therefore, we confirm that the β -naphthylene structure is mechanically stable under small deformations.

One can note that Young's modulus (Y), which quantifies the material's resistance to uniaxial deformation, exhibits strong anisotropy. As shown in Fig. 9a, Y reaches a maximum of 165.88 N/m along one crystallographic direction and a minimum of 43.71 N/m in the perpendicular direction, resulting in an anisotropy ratio of 3.79. This trend suggests that the material's stiffness is strongly dependent on the applied strain direction [61]. The maximum Young's modulus

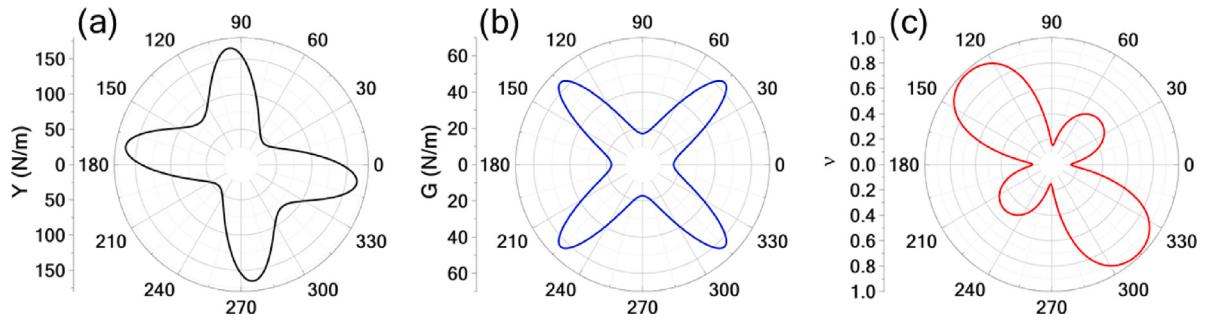


Fig. 9. Polar plots of the in-plane mechanical properties of β -naphthylene: (a) Young's modulus (Y), (b) shear modulus (G), and (c) Poisson's ratio (ν).

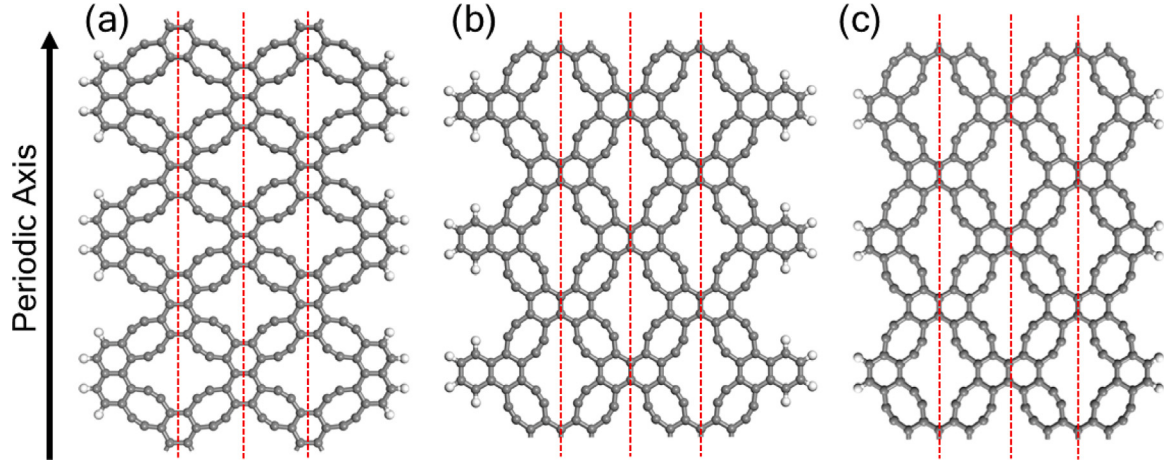


Fig. 10. Schematic representation of the 1D structures of β -naphthylene nanoribbons. (a) Zigzag nanoribbon (ZNNR), (b) armchair nanoribbon with double hexagon edges (ANNR- hh), and (c) armchair nanoribbon with single hexagon edges (ANNR- h).

Table 2

Maximum and minimum values of Young's modulus ($Y_{\max}$, $Y_{\min}$) in N/m, Poisson ratio ($\nu_{\max}$, $\nu_{\min}$), and Shear modulus ($G_{\max}$, $G_{\min}$) in N/m for β -naphthylene and other carbon monolayers. The values reported in this table were calculated in this study.

	$Y_{\max}/Y_{\min}$	$\nu_{\max}/\nu_{\min}$	$G_{\max}/G_{\min}$
β -naphthylene	165.88/43.71	0.98/0.15	63.94/17.17
PHE-graphene	262.29/262.29	0.26/0.26	103.91/103.91
Graphenylene	209.02/209.02	0.27/0.27	82.11/82.11
Graphene	345.42/345.42	0.17/0.17	147.60/147.60
Graphenyldiene	122.47/122.47	0.35/0.35	45.29/45.29
Penta-graphene	271.81/266.67	-0.08/-0.10	151.21/144.98
T-graphene	293.90/148.02	0.16/0.58	126.57/148.02
Graphyne	123.59/123.59	0.45/0.45	42.63/42.63

of β -naphthylene surpasses that of graphenyldiene (122.47 N/m) and graphyne (123.59 N/m) (refer to Table 2).

The shear modulus (G), which describes the resistance to shear deformations, also demonstrates notable directional variation. Fig. 9b depicts that G varies from 63.93 N/m to 17.17 N/m, corresponding to an anisotropy ratio of 3.73. Interestingly, the maximum shear modulus ($G_{\max}$) exceeds that of graphenyldiene (45.29 N/m) and graphyne (42.63 N/m), suggesting that β -naphthylene resists shear forces more effectively than other related 2D carbon materials. The Poisson's ratio (ν), shown in Fig. 9c, exhibits the highest degree of anisotropy, ranging from 0.98 to 0.15, with an anisotropy ratio of 6.48. This extreme variation indicates that β -naphthylene's ability to contract in one direction when stretched in another highly depends on the strain orientation. This wide range of mechanical properties in β -naphthylene is comparable to that of several anisotropic materials ever reported, such as T-graphene [62], Ographene [63], and 2D MoO_3 [64], and Tetrahex Carbon [65]. Furthermore, this feature also highlights its potential

for applications that require directional mechanical responses, such as mechanical metamaterials [66], strain-sensitive output devices [67,68], and flexible electronics [69].

3.6. One-dimensional (1D) structures

We now turn our attention to nanoribbons derived from β -naphthylene. Due to their rectangular lattice, the nanoribbons were generated by exploring the cutting directions of the armchair (y or b) and zigzag (x or a). It is worth noting that experimental control over edge terminations in atomically precise nanoribbons remains a challenge. Techniques such as bottom-up synthesis or post-growth chemical functionalization may be necessary to stabilize the edges of well-defined nanoribbons and preserve their electronic properties [70–72].

The armchair and zigzag β -naphthylene nanoribbons are referred to as ANNRR and ZNNR, respectively. For ZNNR, we considered a single type of edge in the non-periodic direction, consisting of two hexagons, as illustrated in Fig. 10a. In the case of ANNRRs, two edge terminations were examined: one with two hexagons (hh) (Fig. 10b) and another with a single hexagon (h) (Fig. 10c), labeled ANNRR- hh and ANNRR- h , respectively. The chosen nanoribbon width spans three naphthyl groups. This methodology is in agreement with the approach adopted by Paz et al. [32] in their investigation of the 1D structures of β -naphthylene. Note that β -naphthylene nanoribbons did not show reconstructions or bond breaks in optimization. In addition, the planarity of 2D β -naphthylene was maintained.

To explore the electronic properties of the β -naphthylene 1D structures, we calculated the band structures for ZNNR, ANNRR- h , and ANNRR- hh , as depicted in Fig. 11. The ZNNR retains a metallic character, mirroring its 2D counterpart. However, a remarkable characteristic emerges with a Dirac cone positioned closer to the Γ point. In contrast,

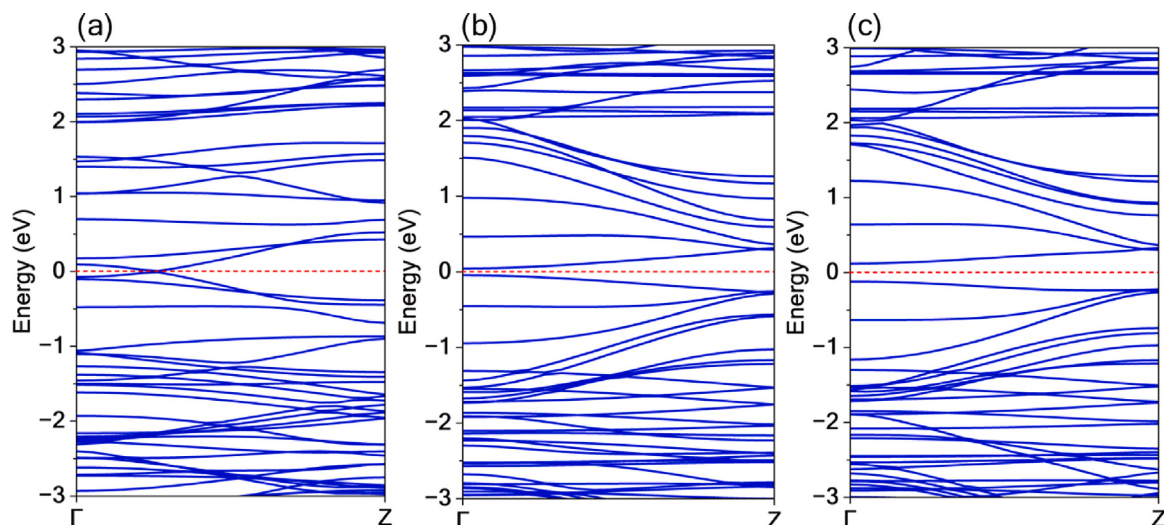


Fig. 11. Electronic band structures of β -naphthylene nanoribbons: (a) zigzag-edge nanoribbon (ZNNR), (b) armchair-edge nanoribbon with hollow-hollow edges (ANNR-*hh*), and (c) armchair-edge nanoribbon with hollow edges (ANNR-*h*); the red dashed line denotes the Fermi level (E_F).

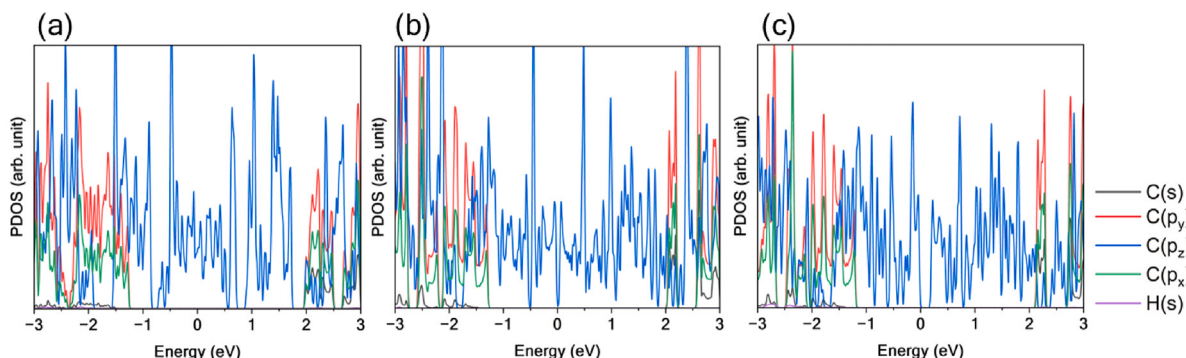


Fig. 12. Projected density of states (PDOS) for β -naphthylene nanoribbons: (a) ZNNR, (b) ANNR-*hh*, and (c) ANNR-*h*.

the ANNR-*hh* and ANNR-*h* structures reveal a transition from metallic to semiconducting behavior, with direct band gaps at the Γ point of 0.08 eV and 0.24 eV, respectively. This band gap opening can be attributed to the edge-induced symmetry-breaking and quantum confinement effects inherent to the reduced dimensionality.

To verify the composition of the electronic states of the β -naphthylene nanoribbons, the projected density of states (PDOS) was calculated, as shown in Fig. 12. The general features closely resemble those of the 2D counterpart, such as the prominent contribution of p_y and p_x states at the lower and higher energies of the valence band (VB) and conduction band (CB), respectively. Around the Fermi level (E_F), the p_z states dominate entirely, with no significant contributions from other orbitals. Furthermore, hydrogen states do not show relevant influence across the evaluated energy range.

For ZNNR (Fig. 12a), despite the presence of a Dirac cone, the PDOS reveals a non-zero density of states at the Fermi level. This finite DOS arises from the low dispersion of bands near the Dirac point, causing an accumulation of electronic states around E_F , preventing the DOS from vanishing completely. In the cases of ANNR-*hh* (Fig. 12b) and ANNR-*h* (Fig. 12c), a noticeable accumulation of states at the band edges is observed, accompanied by an abrupt drop in the DOS at these regions, characteristic of highly correlated electronic states.

4. Conclusion

In this work, we explore the structural, electronic, mechanical and optical properties of a novel porous 2D carbon allotrope, β -naphthylene,

using first-principles calculations. Our results confirmed its stability, with phonon dispersion and *ab initio* molecular dynamics (AIMD) simulations indicating a robust framework.

The electronic band structure revealed that β -naphthylene exhibits metallic behavior, with strong π -character for the states in the band edges. Mechanical analysis highlighted its strong anisotropy, showing a significant directional dependence on Young's modulus, shear modulus, and Poisson's ratio. This mechanical anisotropy surpasses that of previously studied 2D carbon materials, positioning β -naphthylene as a promising candidate for strain engineering applications.

Optical investigations demonstrated that β -naphthylene exhibits strong IR absorption, moderate interaction within the visible spectrum, and pronounced UV absorption peaks, suggesting its potential for optoelectronic and photonic technologies. Furthermore, the Raman and IR spectra provided distinctive vibrational fingerprints, providing a reliable reference for future experimental validation.

We further examine the electronic properties of 1D β -naphthylene nanoribbons. Although ZNNR retained the metallic character of its 2D counterpart, ANNR-*h* and ANNR-*hh* showed a transition from metallic to semiconducting behavior, with direct band gaps of 0.08 and 0.24 eV, respectively. These findings underscore the tunable electronic properties of β -naphthylene, driven by dimensionality reduction and edge geometry, broadening its potential for nanoelectronic applications.

CRedit authorship contribution statement

José A.S. Laranjeira: Writing – review & editing, Writing – original draft, Methodology, Formal analysis, Conceptualization. **K.A.L. Lima:**

Writing – review & editing, Writing – original draft, Methodology, Investigation, Formal analysis, Conceptualization. **Nicolas F. Martins:** Writing – review & editing, Writing – original draft, Methodology, Investigation, Formal analysis, Conceptualization. **Marcelo L.P. Junior:** Writing – review & editing, Writing – original draft, Methodology, Formal analysis, Conceptualization. **L.A. Ribeiro Junior:** Writing – review & editing, Writing – original draft, Methodology, Formal analysis, Conceptualization. **Julio R. Sambrano:** Writing – review & editing, Writing – original draft, Methodology, Formal analysis, Conceptualization.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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Data availability

Data will be made available on request.

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