



On the structural, electronic, and optical properties of L-histidine crystal: a DFT study

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Abstract

Context The monoclinic L-histidine crystal is critical for protein structure and function and is also found in the myelin of brain nerve cells. This study numerically examines its structural, electronic, and optical properties. Our findings indicate that the L-histidine crystal has an insulating band gap of approximately 4.38 eV. Additionally, electron and hole effective masses range between $3.92m_0$ – $15.33m_0$ and $4.16m_0$ – $7.53m_0$, respectively. Furthermore, our investigation suggests that the L-histidine crystal is an excellent UV collector due to its strong optical absorption activity for photon energies exceeding 3.5 eV.

Methods To investigate the structural, electronic, and optical properties of L-histidine crystals, we used the Biovia Materials Studio software to conduct Density Functional Theory (DFT) simulations as implemented in the CASTEP code. Our DFT calculations were performed using the generalized gradient approximation (GGA) as parameterized by the Perdew-Burke-Ernzerhof (PBE) exchange-correlation functional, with an additional dispersion energy correction (PBE + TS) based on the model proposed by Tkatchenko and Scheffler to describe van der Waals interactions. Additionally, we employed the norm-conserving pseudopotential to treat core electrons.

Keywords L-Histidine crystal · Electronic structure · Optical properties · Structural properties · DFT calculations

Introduction

Histidine is an essential amino acid with an imidazole functional group in its side chain, giving it aromatic properties [49]. It is one of the twenty standard amino acids found in proteins in all living organisms, existing in both D-histidine and L-histidine isomeric forms [50]. This amino acid plays a significant role in various biological mechanisms, including the formation of hemoglobin [51], and is also used to treat conditions such as allergic diseases and anemia [1–4].

One critical feature of Histidine is that it's the only amino acid whose side chain can shift from a non-protonated to a protonated state, allowing it to exhibit both acid and nucleophilic functionality [52]. The non-protonated state plays a

crucial role as a base for the conformational stability of proteins [1, 5]. Additionally, the structural stability of amino acid crystals is currently an active area of research, as these structures have the potential for application in technological devices [3].

L-Histidine is a noteworthy amino acid as the human body cannot synthesize it and must be obtained through diet [7–9]. It plays a crucial role in many biological functions, such as the development and maintenance of myelin, which is responsible for communication between the brain and body, as well as the regulation of essential metal levels like zinc, iron, copper, and nickel [10]. L-Histidine has many applications, including biosynthesis [1, 4], metabolism, physiological importance, and supplementation [11]. It is also involved in human genome transcription factors [12], L-histidine capped Ag and Au nanoparticles for dopamine detection [13], and protein interactions [14].

Recently, L-histidine has attracted significant interest for both experimental and theoretical researchers [53–59]. Monocrystalline samples of L-histidine hydrochloride monohydrate were analyzed using Raman spectroscopy to

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investigate the effect of temperature changes on the vibrational spectrum [3]. Y. Wang and colleagues developed a straightforward, non-destructive method using terahertz frequency-domain spectroscopy to quickly and accurately analyze L-histidine and α -lactose in dietary supplements [6]. Their findings demonstrated high accuracy levels ranging from 94.8 to 110% for L-histidine and 98.9 to 110% for α -lactose, comparable to those of ion chromatography and high-performance liquid chromatography, respectively [6]. These results highlight the potential of the developed method for rapid, in situ analysis of nutrients in dietary supplements.

To compare theoretical and experimental results, terahertz spectroscopy and density functional theory (DFT) calculations were used to analyze the low-frequency vibration of pure solid-state L-histidine crystals, and their monohydrochloride monohydrate [15]. Results showed notable differences in the contribution of intermolecular movement and the distortions of the dihedral angle of the imidazole ring between the crystals. In another study, the spectroscopic properties of the stable semi-organic complex L-histidine tetrafluoroborate were also investigated [16]. It was demonstrated that Infra-Red and Raman spectroscopy's experimental and DFT calculation data agree well. The findings also indicated that L-histidine tetrafluoroborate is a stable complex suitable as a new violet-light emitting material with high non-linear properties and potential applications as an efficient optical limiter.

Previous studies have shown that molecules and crystals from the L-histidine family exhibit non-linear optical effects due to the presence of the imidazole ring [16–22]. This property is common to other amino acid crystals [23–28], suggesting that these structures are potential candidates for second harmonic generators, frequency converters, light modulators, optical commutators, and optical memory storage applications. The unique features and significant role of L-histidine crystals discussed earlier suggest that a comprehensive investigation into their properties could pave the way for their use in diverse optoelectronic applications.

The present study used DFT calculations to investigate L-histidine crystals' structural, electronic, and optical properties. Our findings suggest that the crystal has an insulating band gap of approximately 4.38 eV, with electron and hole effective masses ranging between $3.92m_0$ – $15.33m_0$ and $4.16m_0$ – $7.53m_0$, respectively. Our investigation also suggests that L-histidine crystals are excellent UV collectors, displaying strong optical absorption for photon energies exceeding 3.5 eV. This paper is organized as follows: the “Methodology” section describes the computational approach used to obtain the properties of interest, while “Structural properties”, “Electronic properties”, and “Optical properties” sections discuss the results for the structural, electronic, and optical properties of model L-histidine crystals.

Finally, the “Conclusions” section presents the concluding remarks.

Methodology

The monoclinic form of the L-histidine crystal, with a chemical formula of $C_6H_9N_4O_2$, was studied. This amino acid is composed of four functional groups: carboxyl (COO^-), amino (NH_3^+), imidazole ring ($C_3N_2H_3$), and lateral chain ($C_5N_2H_7$), as illustrated in Fig. 1a). The unit cell of the monoclinic L-histidine (mLH) crystal, consisting of two molecules and their respective hydrogen bonds, is shown in Fig. 1b). In Fig. 1c), a $(3 \times 1 \times 3)$ L-histidine crystal supercell is displayed. The Biovia Materials Studio software was used to conduct DFT simulations as implemented in the CASTEP code [36, 37]. We used the generalized gradient approximation (GGA) as parameterized by the Perdew-Burke-Ernzerhof (PBE) exchange-correlation functional [35], with an additional dispersion energy correction (PBE + TS) based on the model proposed by Tkatchenko and Scheffler [39] to describe van der Waals interactions. We employed norm-conserving pseudopotentials to treat core electrons [41].

Experimental data indicate that the mLH crystal forms four distinct hydrogen bonds: $N3-H \cdots O1$ (at $x, y, -1 + z$), $N1-H \cdots O2$ (at $-1 + x, y, z$), $N1-H \cdots O2$ (at $1 - x, \frac{1}{2} + y, -1 + z$), and $N1-H \cdots N2$ (intramolecular) [40]. The molecule's conformation is stabilized by the intramolecular hydrogen bond between $N1-H \cdots N2$ [40].

The unit cell of the mLH crystal was geometrically optimized using the following set of converge criteria for the self-consistent field (SCF) calculations: SCF tolerance of 5×10^{-6} eV/atom, a maximum force on each atom of 0.01 eV/Å, pressure limit of 0.01 GPa, and maximum atomic displacement of 5×10^{-4} Å, and cutoff energy of about 830 eV. To optimize the unit cell, we employed the BFGS (Broyden, Fletcher, Goldfarb, and Shanno) algorithm [42]. The tolerance for each SCF step was tested to achieve convergence, ensuring that the total/atom energy variation is approximately 5×10^{-7} eV/atom and the eigenvalues are less than 3.33×10^{-7} eV for three successive steps.

Results

Structural properties

The mLH crystal, as classified by X-ray experimental data, belongs to the spatial group $P2_1$ and has the following crystal parameters (refer to Fig. 1): $a=5.17$ Å, $b=7.38$ Å, $c=9.49$ Å, $\beta = 97.16^\circ$, and $V=359.00$ Å³ [40]. Our theoretical calculations, which are in good agreement with the experimental

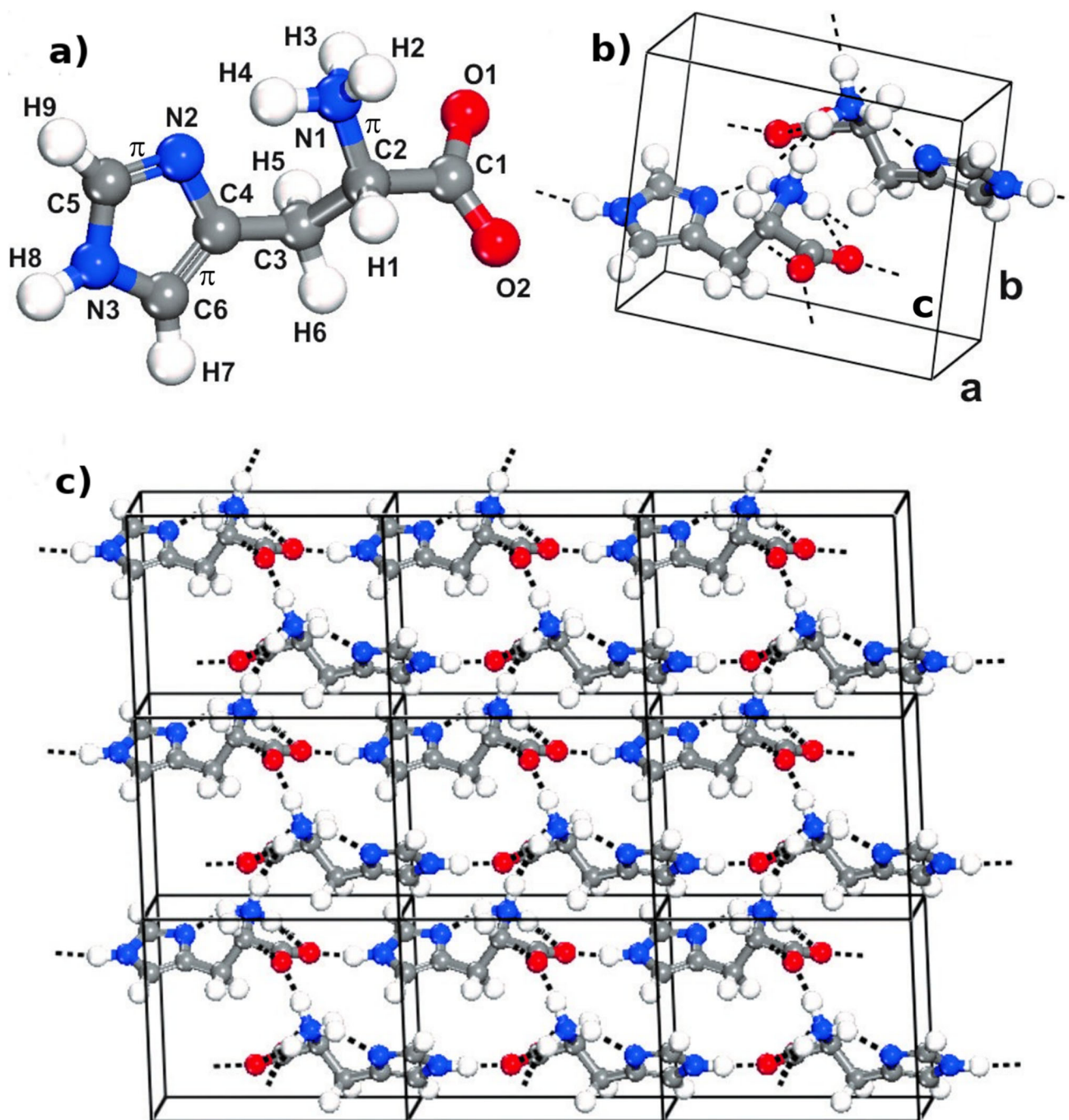


Fig. 1 Panel (a) depicts the schematic representation of the LH molecule with its constituent atoms represented as follows: oxygen (red), nitrogen (blue), carbon (gray), and hydrogen (white). Panel (b)

displays the unit cell of the mLH crystal, while panel (c) shows a $(3 \times 1 \times 3)$ mLH crystal supercell, where the hydrogen bonds are highlighted

data, yielded: $a=5.13 \text{ \AA}$, $b=7.36 \text{ \AA}$, $c=9.47 \text{ \AA}$, $\beta = 97.76^\circ$, $V=354.80 \text{ \AA}^3$.

When comparing the atomic positions obtained from our theoretical calculations to experimental data, we found that the x-coordinate of H1 bonded to C2 had the largest deviation, with $\Delta x = -0.091 \text{ \AA}$. The y- and z-coordinates of the H2 and

H4 hydrogen atoms bonded to N1 also showed deviations, with $\Delta y = -0.068 \text{ \AA}$ and $\Delta z = -0.049 \text{ \AA}$, respectively. We observed deviations in the theoretical hydrogen bond lengths, with the C1-C2 bond having the largest relative deviation at about 23.323%, while the others did not exceed 3.624%. Additionally, we obtained a relative deviation of about 7.4%

in the angle of the C3-C4-N2 bond compared to experimental data, which is the largest one. These results suggest that our approach accurately reproduces the structural parameters reported in the experiment [40] and is suitable for analyzing mLH crystals' electronic and optical properties.

Electronic properties

We now discuss the electronic properties of the mLH crystal. Table 1 shows the electrical charge per atom, and functional group, distributions in the L-histidine molecule obtained from Mulliken's and Hirshfeld's electronic population analysis [43–45]. The results are expressed in elementary charge unit e . Our analysis reveals that the C1 has the highest positive charge, with Mulliken's (Hirshfeld's) charges of about $0.66e$ ($0.16e$), while the N1 has the highest negative charge, with Mulliken's (Hirshfeld's) charges of about $-0.77e$ ($-0.07e$).

Table 2 presents Mulliken's and Hirshfeld's electronic population for the carboxyl (COO^-), amino, imidazole ring, and lateral chain functional groups. The corresponding Mulliken's (Hirshfeld's) charges are about $-0.67e$ ($-0.31e$), $0.09e$ ($0.06e$), $0.23e$ ($0.10e$), and $1.59e$ ($0.23e$), respectively. These findings indicate a dipolar behavior of the zwitterionic state inside the crystal [45–47]. It is worth noting that the zwitterionic structure of L-histidine holds, even in the gas phase, as reported in references [60, 61]. As a result, the atomic charge configuration tends to be also similar.

Table 1 Mulliken and Hirshfeld electronic population analysis of mLH crystal

Atoms	Mulliken	Hirshfeld
H1	+0.28	+0.04
H2	+0.29	+0.04
H3	+0.27	+0.04
H4	+0.30	+0.05
H5	+0.28	+0.04
H6	+0.43	+0.10
H7	+0.41	+0.10
H8	+0.43	+0.10
H9	+0.43	+0.10
C1	+0.66	+0.16
C2	-0.15	+0.02
C3	-0.52	-0.07
C4	+0.08	+0.01
C5	+0.06	+0.04
C6	-0.16	-0.03
N1	-0.77	-0.07
N2	-0.51	-0.16
N3	-0.51	-0.06
O1	-0.66	-0.24
O2	-0.67	-0.23

Table 2 Mulliken and Hirshfeld electronic population analysis of carboxyl, amino, imidazole ring and lateral chain functional group

Atoms	Mulliken	Hirshfeld
Carboxyl	-0.67	-0.31
Amino	+0.09	+0.06
Imidazole ring	+0.23	+0.10
Lateral chain	+1.59	+0.23

Figure 2 illustrates the high symmetric k-point path in the first Brillouin zone. The high-symmetry k-points are Γ (0,0,0), Z (0,0,1/2), Y (0,1/2,0), A (-1/2,1/2,0), B (-1/2,0,0), C (0,1/2,1/2), D (-1/2,0,1/2), and E (-1/2,1/2,1/2).

The band structure configuration and the corresponding total density of states (DOS) of mLH are shown in Fig. 3. The top panels, i.e., Fig. 3a, b), and c), provide an overview

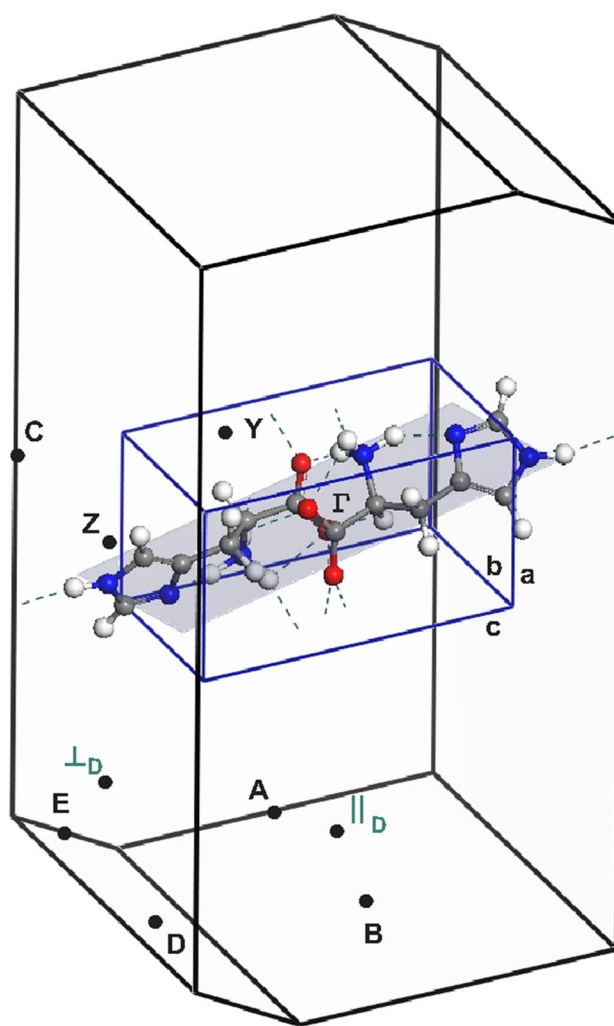


Fig. 2 The first Brillouin zone of the mLH crystal, including the high-symmetry k -points of the reciprocal lattice (designated with capital letters)

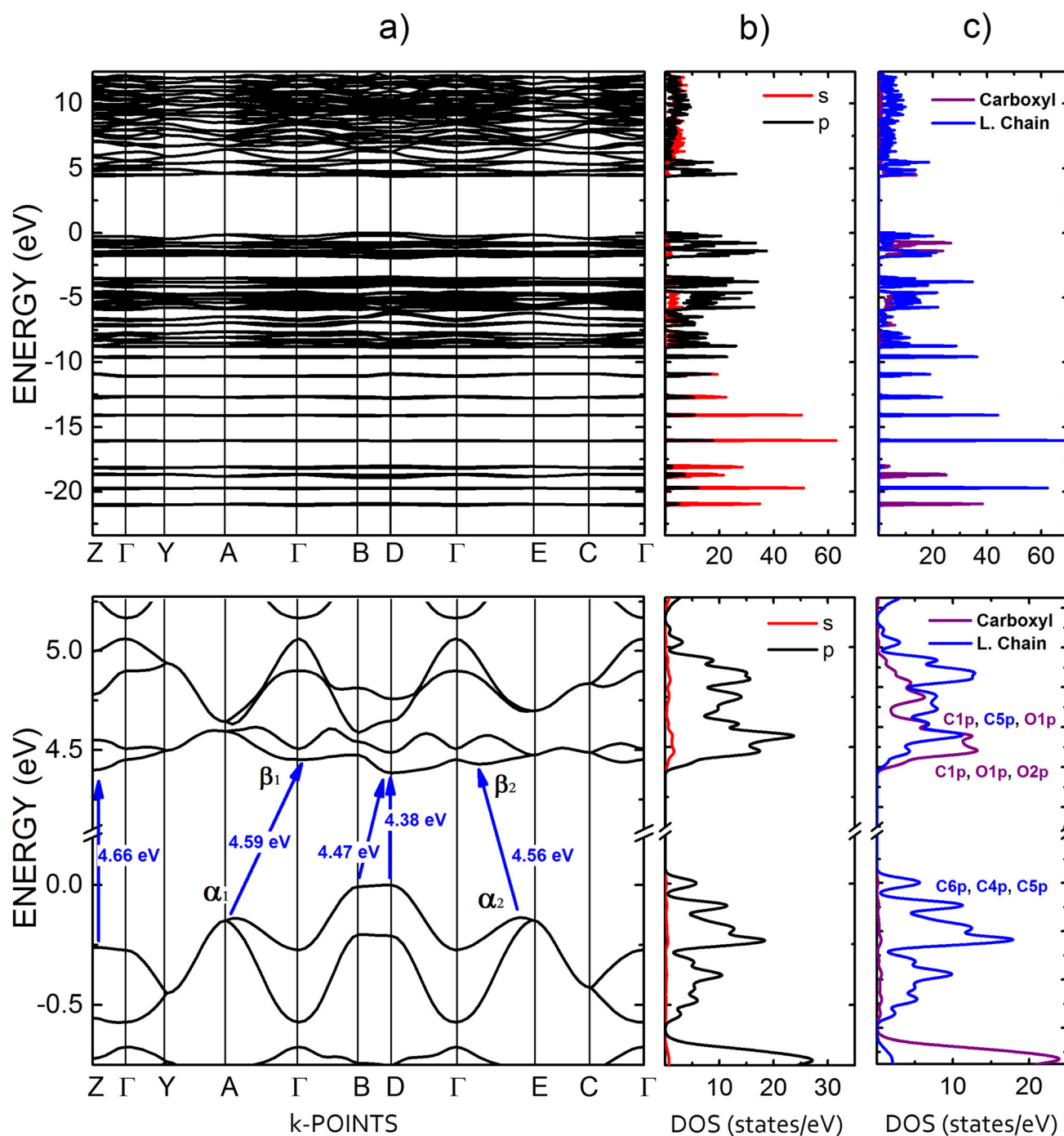


Fig. 3 Band structure configuration, and related total density of states (DOS), of mLH. Panels a), b), and c) depict the full bandgap structure, the DOS for the mLH crystal, and the DOS for the carboxyl and lateral

chain functional groups, respectively. The bottom panels offer a closer look at these quantities for energies ranging from -0.7 eV to 5.2 eV. The energy levels were shifted according to the Fermi level

of the full bandgap structure, the DOS for the mLH crystal, and the DOS for the carboxyl and lateral chain functional groups, respectively. The bottom panels offer a closer look at these quantities for energies ranging from -0.7 eV to 5.2 eV, with the energy levels shifted relative to the Fermi level. Notably, the mLH crystal exhibits an insulating direct band gap of approximately 4.38 eV along the $D \rightarrow D$ path. Other possible indirect electron transfer channels include the $B \rightarrow D$ (4.47 eV), $A \rightarrow \Gamma$ (4.59 eV, labeled as $\alpha_1 \rightarrow \beta_1$), $E \rightarrow \Gamma$ (4.56 eV, labeled as $\alpha_2 \rightarrow \beta_2$), and $Z \rightarrow Z$ (4.66 eV) paths.

Figure 3b) (top panel) shows that in the lowermost valence band, the s orbital dominates, while s and p orbitals contribute to form the uppermost level of the conduction band. In Fig. 3b) (bottom panel), we see that the p orbitals are prominent at the top of the valence band as well as in the lowermost energy levels of the conduction band. Lastly, Fig. 3c) (top panel) shows the DOS for the carboxyl and lateral chain, where the lateral chain is dominant, with the prevalence of the p orbital for both valence and conduction bands. Addi-

tionally, in the valence band, the major contribution is from the linear chain through the C4, C5, and C6, while in the conduction band, there is a mixed contribution of both the carboxyl group through the C1, O1, and O2 atoms and the lateral chain group with the C5, as shown in Fig. 3c) (bottom panel).

We further analyze the electronic results by presenting the effective mass (m^*) of electrons (m_e) in the conduction band and holes (m_h) in the valence band. The inverse of the second derivative of the energy (E) concerning distance in reciprocal space (k) is proportional to m_e and m_h , respectively. The following equation can express these values [45, 48]:

$$\frac{1}{m^*} = \frac{1}{\hbar^2} \frac{\partial^2 E}{\partial k^2}, \quad (1)$$

where $\hbar$ is the reduced Planck's constant. The band curvature at points of interest in the Brillouin zone was calculated for the mLH crystal by fitting quadratic curves, as shown in

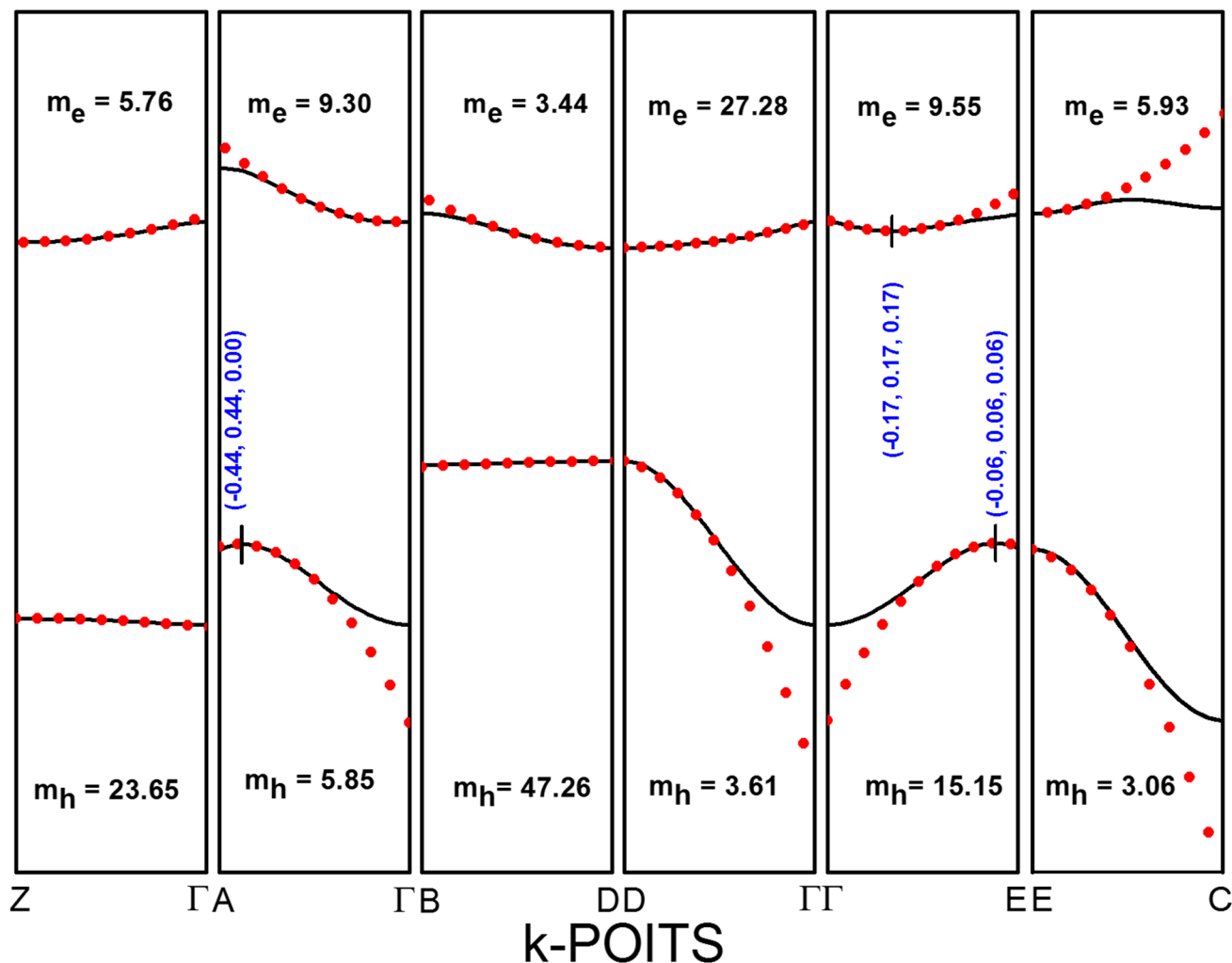


Fig. 4 Energy level fittings to estimate the electron and hole effective masses in mLH crystals

Fig. 4. The effective mass is an important factor in determining the viability of charge transport through the valence and conduction bands. In this way, in Fig. 4, it can be seen that the hole effective masses vary from $3.06m_0$ (at the $E \rightarrow C$ path) to $47.26m_0$ (at the $B \rightarrow D$ path), while the electron masses vary from $3.44m_0$ (at the $B \rightarrow D$ path) to $27.28m_0$ (at the $D \rightarrow \Gamma$ path), where m_0 represents the electron mass.

Figure 5 highlights the valence and conduction bands for special directions, i.e., perpendicular and parallel to the molecular plane of the mLH crystal, indicated by $\perp_D$ and $\parallel_D$ (refer to Fig. 2). These directions represent the maximum point of the valence band and the minimum point of the conduction band, shown in Fig. 5a) and b) (top panels), respectively. Their related valence and conduction bandwidths (bottom panels) allow evaluation of the effective mass character, with a narrow bandwidth corresponding to a large effective mass and a wide bandwidth to a small effective mass.

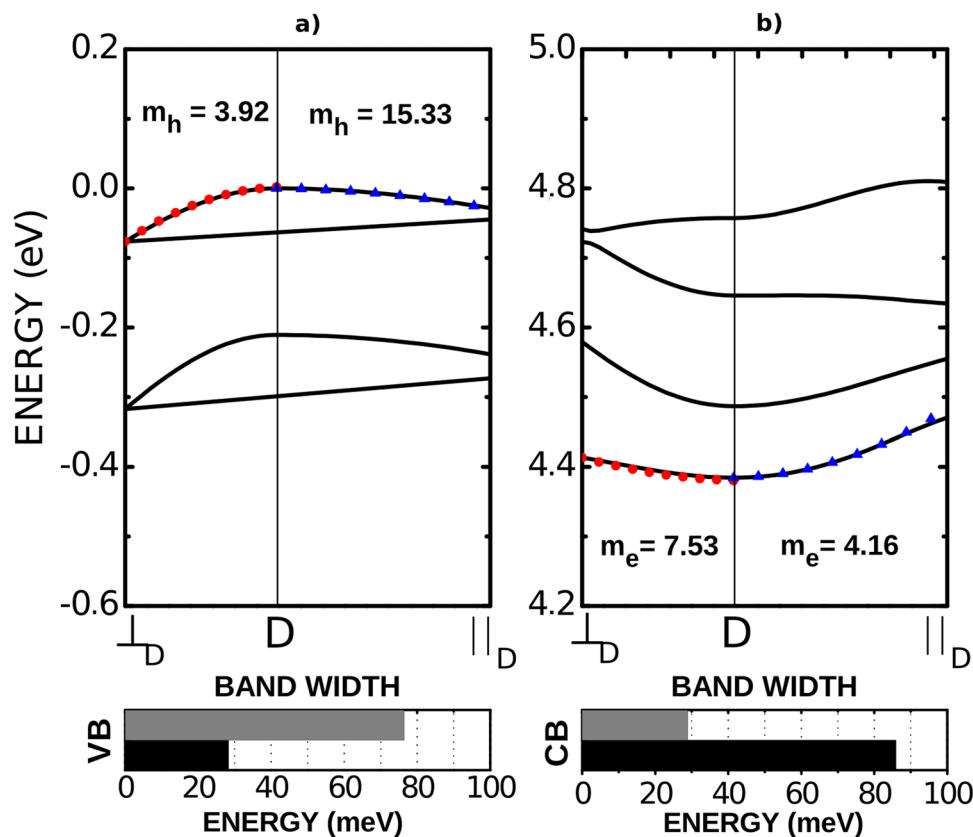
In Fig. 5a) (bottom panel), the valence bandwidth reveals that the effective mass for the $\perp_D$ direction (gray bar) is larger than the one for the $\parallel_D$ direction (black bar). On the other hand, Fig. 5b) (bottom panel) shows that the effective mass for the $\perp_D$ direction (gray bar) is lower than the one for the $\parallel_D$ direction (black bar) in the conduction band. Moreover, the valence band $\perp_D$ direction has an energy of about 77 meV with $m_h = 3.93m_0$, whereas the $\parallel_D$ direction presents

an energy of about 28 meV with $m_h = 15.33m_0$. Regarding the conduction band and $\perp_D$, the energy is about 29 meV with $m_e = 7.53m_0$, while for the $\parallel_D$ direction, the energy is about 78 meV with $m_e = 4.16m_0$. It is worth emphasizing that the directions with the lowest effective mass values are more favorable for electronic transference.

Now, we analyzed the partial density of states (PDOS) per atom and functional group. Figure 6a) shows the relative contributions of each atomic species C, O, N, and H. We found that the C, O, and N atoms contribute significantly with p orbitals close to the bandgap for both the valence and conduction bands, whereas the H atom contributes only with s orbitals. Figure 6b) displays the electronic contribution of the carboxyl, amino, imidazole ring, and lateral chain functional groups. We observed that s and p orbitals contribute to both functional groups, with the p orbitals being more prevalent near the valence and conduction bands for the carboxyl, imidazole ring, and lateral chain.

The s orbital contributes prominently only to the most negative energy range of the conduction band in the energy range from -21.25 eV to -11.25 eV for all functional groups constituting the mLH crystal. The amino group has the lowest contribution at the valence and conduction bands. For this group, the s and p orbitals overlap in the vicinity of both bands with PDOS peaks at about 0; electron/eV. On the other hand, the imidazole ring and lateral chain groups

Fig. 5 Top panels show the (a) valence and (b) conduction bands for special directions, i.e., perpendicular and parallel to the molecular plane of the mLH crystal, indicated by $\perp_D$ and $\parallel_D$ (refer to Fig. 2). The bottom panels depict the related valence and conduction bandwidths



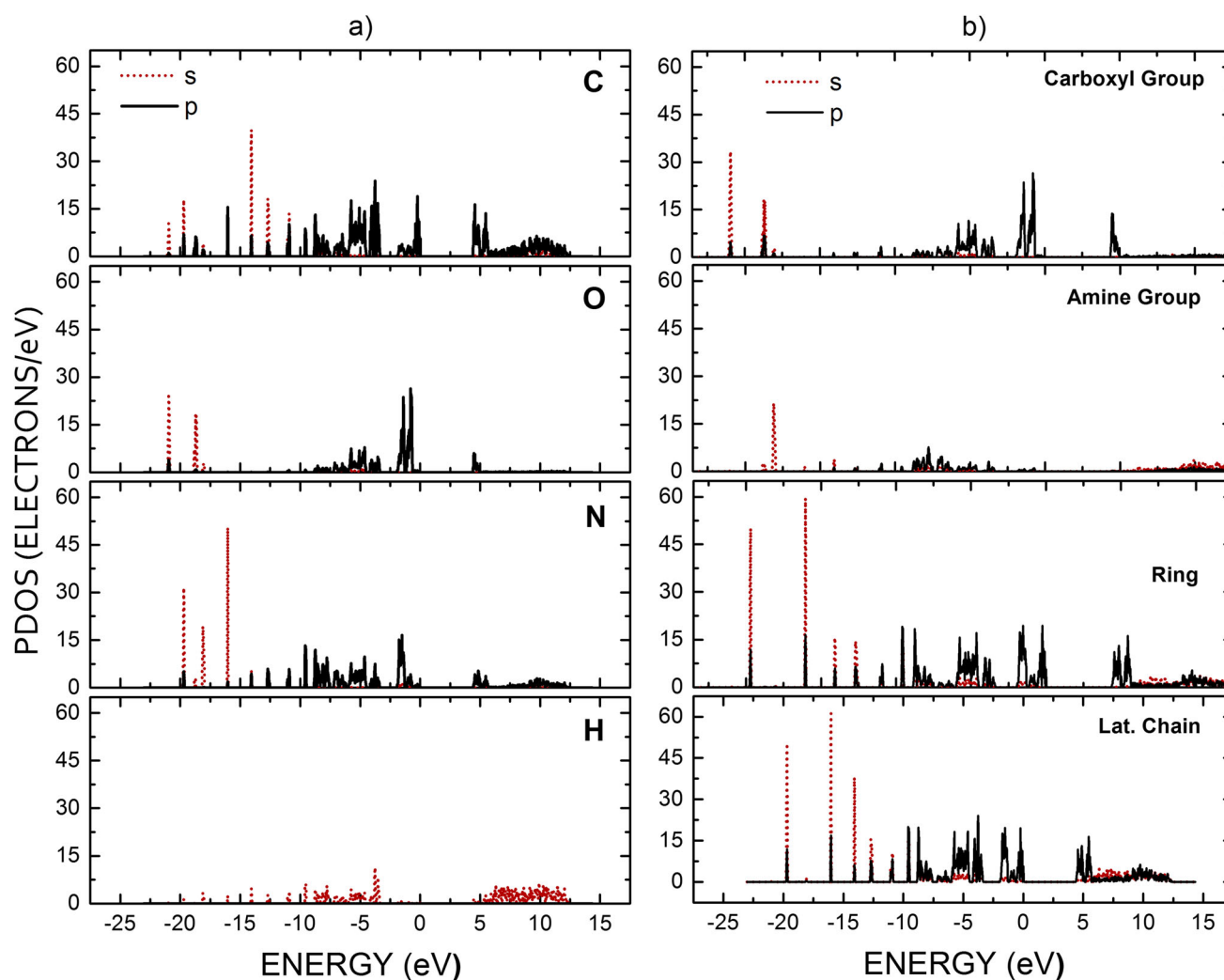


Fig. 6 Partial density of states (PDOS) a) per atom and b) functional groups

present prominent peaks at the valence and conduction bands, differing from the carboxyl group in the same region. This PDOS analysis sheds light on the origin of the relevance of the band structure features of the mLH crystal.

Optical properties

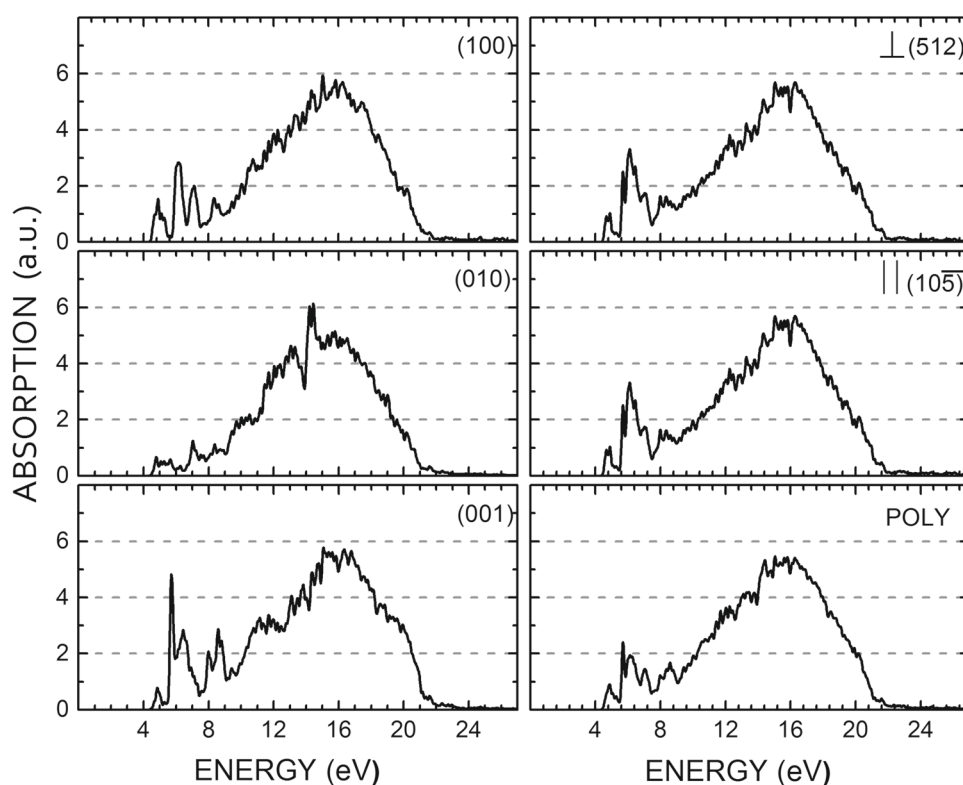
Here, we discuss the optical properties of the mLH crystals. We consider the incident light polarized along the (001), (010), (100), (512) and $(10\bar{5})$ directions, as well as one polycrystalline sample. The incident light planes (512) and $(10\bar{5})$ correspond to the $\perp_D$ and $\parallel_D$ special directions (refer to Fig. 2).

Figure 7 shows the optical absorption behavior of the mLH crystal when the incident light is polarized along the selected crystal planes mentioned earlier. The absorp-

tion spectrum varies in the energy range of approximately 4.0 eV to 24.0 eV, indicating an intense optical absorption activity in the ultraviolet region (UV). This absorption range is important for optoelectronic devices that act as UV collectors.

Notably, a prominent peak of light absorption is observed in the (001), (100), (512), and $(10\bar{5})$ directions at about 4.0 eV to 10.0 eV, with emphasis on the (001) direction. The (010) direction shows low optical absorption activity in the same range. The energy range from about 10.0 eV to 24.0 eV shows the region of the greatest absorption, with peak absorption around 16.0 eV. The mLH crystal exhibits anisotropic behavior, except for the (512) and $(10\bar{5})$ directions. Moreover, the polycrystalline sample shows the typical overall behavior of optical absorption seen in the (001), (100), (512), and $(10\bar{5})$ directions.

Fig. 7 Optical absorption of mLH crystal along the (001), (010), (100), (512) and (10 $\bar{5}$) incidence directions. The optical absorption for a polycrystalline sample is also shown



Finally, Fig. 8 displays the real (black) and imaginary (red) parts of the dielectric function of the photon energy. The dielectric function exhibits an anisotropic behavior along the incident planes, except for the (001) and (10 $\bar{5}$) directions, which have the same behavior. The (001) and (10 $\bar{5}$) directions have the most significant peaks in terms of amplitude for the (001), (100), (512), and (10 $\bar{5}$) directions. The (010) direction has no peak in the same region. Furthermore, the amplitude of the real and imaginary parts of the dielectric function varies, indicating the presence of a strong dipole moment induced by the electric field of the incident light. The dielectric function of the polycrystalline sample shows the typical overall behavior for (001), (100), (512), and (10 $\bar{5}$) directions.

The amplitude of the real and imaginary parts of the dielectric function varies due to a strong dipole moment induced by the light [62]. Light interacting with material creates a polarization effect that moves electrons and generates a dipole moment [63]. This dipole moment interacts with the light's electromagnetic field and is reflected in the dielectric function [63]. The real part of the dielectric function represents the material's ability to store electrical energy, while the imaginary part represents its ability to dissipate electrical energy [64]. The variation in the amplitude of these two parts shows that the material responds differently to the electric field of the light due to the dipole moment induced by the light [64]. Thus, the strength of the induced dipole moment affects the

amplitude of the real and imaginary parts of the dielectric function.

Conclusions

In summary, we have investigated the optical properties of mLH using first-principles calculations. Our results showed that our calculated structural properties matched well with experimental findings. The electronic band structure revealed that mLH is an insulator material with a direct band gap of about 4.38 eV. Furthermore, effective mass calculations indicated the possibility of charge transport through valence and conduction bands, especially along directions with the lowest effective mass values. The hole effective masses vary from $3.06m_0$ to $47.26m_0$, while the electron masses vary from $3.44m_0$ to $27.28m_0$.

The mLH crystal exhibits anisotropic behavior in optical absorption and dielectric function along different incident planes. The absorption spectrum varies in the energy range of approximately 4.0 eV to 24.0 eV, indicating an intense optical absorption activity in the ultraviolet region (UV). This absorption range is important for optoelectronic devices that act as UV collectors.

Generally, it is difficult to directly compare the optical response of L-histidine with other dyes or materials due to the differences in their molecular structures and electronic prop-

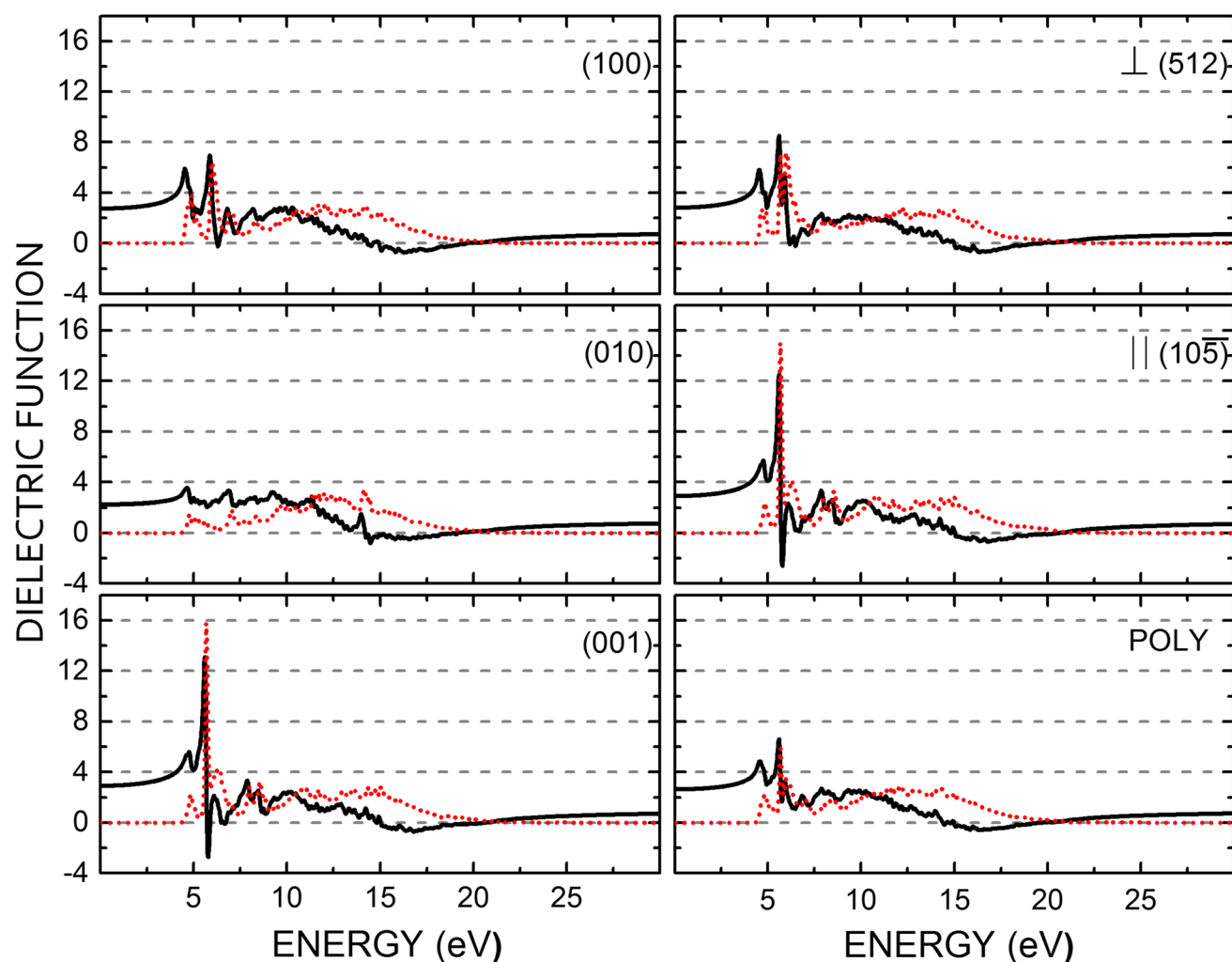


Fig. 8 Real (black) and imaginary (red) parts of the dielectric function for mLH crystal along the (001), (010), (100), (512) and (105) incidence directions. The same quantities are also illustrated for the polycrystalline

erties. The relevance of the optical response of L-histidine depends on the specific application and the device's requirements. For example, if a device needs to detect UV light in a specific range, the absorption spectrum of L-histidine in that range would be compared to that of other materials to determine its effectiveness. L-Histidine's cost, stability, and availability should also be considered when its performance is compared to other materials.

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Author contribution F.A.R.P., A.M.S., N.F.F., R.G.S., and K.A.L.L.: data curation, formal analysis, methodology, and writing — original draft preparation. J.J.S.M., M.L.P.J., and A.M.F.: data curation, formal analysis, methodology, and funding acquisition. L.A.R.J. and V.N.F.: conceptualization, funding acquisition, and writing — reviewing and editing. All authors reviewed the manuscript.

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Code availability N/A

Declarations

Conflict of interest The authors declare no competing interests.

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